

Transcendence in Positive Characteristic

Galois Group Examples and Applications

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Outline

- 1 Preliminaries
- 2 Carlitz logarithms
- 3 Carlitz zeta values
- 4 Rank 2 Drinfeld modules

Preliminaries

- Notation
- Transcendence degree theorem

Scalar quantities

Let p be a fixed prime; q a fixed power of p .

$$A := \mathbb{F}_q[\theta] \quad \longleftrightarrow \quad \mathbb{Z}$$

$$k := \mathbb{F}_q(\theta) \quad \longleftrightarrow \quad \mathbb{Q}$$

$$\bar{k} \quad \longleftrightarrow \quad \overline{\mathbb{Q}}$$

$$k_\infty := \mathbb{F}_q((1/\theta)) \quad \longleftrightarrow \quad \mathbb{R}$$

$$\mathbb{C}_\infty := \widehat{k_\infty} \quad \longleftrightarrow \quad \mathbb{C}$$

$$|f|_\infty = q^{\deg f} \quad \longleftrightarrow \quad |\cdot|$$

Functions

- Rational functions:

$$\mathbb{F}_q(t), \quad \bar{k}(t), \quad \mathbb{C}_\infty(t).$$

- Analytic functions:

$$\mathbb{T} := \left\{ \sum_{i \geq 0} a_i t^i \in \mathbb{C}_\infty[[t]] \mid |a_i|_\infty \rightarrow 0 \right\}.$$

and

$$\mathbb{L} := \text{fraction field of } \mathbb{T}.$$

- Entire functions:

$$\mathbb{E} := \left\{ \sum_{i \geq 0} a_i t^i \in \mathbb{C}_\infty[[t]] \mid \begin{array}{l} \sqrt[i]{|a_i|_\infty} \rightarrow 0, \\ [k_\infty(a_0, a_1, a_2, \dots) : k_\infty] < \infty \end{array} \right\}.$$

Galois groups and transcendence degree

Theorem (P. 2008)

Let M be a t -motive, and let Γ_M be its associated group via Tannakian duality. Suppose that $\Phi \in \mathrm{GL}_r(\bar{k}(t)) \cap \mathrm{Mat}_r(\bar{k}[t])$ represents multiplication by σ on M and that $\det \Phi = c(t - \theta)^s$, $c \in \bar{k}^\times$. Let Ψ be a rigid analytic trivialization of Φ in $\mathrm{GL}_r(\mathbb{T}) \cap \mathrm{Mat}_r(\mathbb{E})$. That is,

$$\Psi^{(-1)} = \Phi \Psi.$$

Finally let

$$L = \bar{k}(\Psi(\theta)) \subseteq \bar{k}_\infty.$$

Then

$$\mathrm{tr. \ deg}_{\bar{k}} L = \dim \Gamma_M \quad (= \dim \Gamma_\Psi).$$

Carlitz logarithms

- Difference equations for Carlitz logarithms
- Calculation of the Galois group
- Algebraic independence
- An explicit example: $\log_C(\zeta_\theta)$

Carlitz logarithms

- Recall the Carlitz exponential:

$$\exp_C(z) = z + \sum_{i=1}^{\infty} \frac{z^{q^i}}{(\theta^{q^i} - \theta)(\theta^{q^i} - \theta^q) \cdots (\theta^{q^i} - \theta^{q^{i-1}})}.$$

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- Its formal inverse is the Carlitz logarithm,

$$\log_C(z) = z + \sum_{i=1}^{\infty} \frac{z^{q^i}}{(\theta - \theta^q)(\theta - \theta^{q^2}) \cdots (\theta - \theta^{q^i})}.$$

- $\log_C(z)$ converges for $|z|_{\infty} < |\theta|^{q/(q-1)}$ and satisfies

$$\theta \log_C(z) = \log_C(\theta z) + \log_C(z^q).$$

The function $L_\alpha(t)$

- For $\alpha \in \bar{k}$, $|\alpha|_\infty < |\theta|^{q/(q-1)}$, we define

$$L_\alpha(t) = \alpha + \sum_{i=1}^{\infty} \frac{\alpha^{q^i}}{(t - \theta^q)(t - \theta^{q^2}) \cdots (t - \theta^{q^i})} \in \mathbb{T},$$

- Connection with Carlitz logarithms:

$$L_\alpha(\theta) = \log_C(\alpha).$$

- Functional equation:

$$L_\alpha^{(-1)} = \alpha^{(-1)} + \frac{L_\alpha}{t - \theta}.$$

Difference equations for Carlitz logarithms

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$$\Phi = \begin{bmatrix} t - \theta & 0 & \cdots & 0 \\ \alpha_1^{(-1)}(t - \theta) & 1 & \cdots & 0 \\ \vdots & \vdots & \ddots & \vdots \\ \alpha_r^{(-1)}(t - \theta) & 0 & \cdots & 1 \end{bmatrix},$$

then Φ represents multiplication by σ on a t -motive M with

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then Φ represents multiplication by σ on a t -motive M with

$$0 \rightarrow \mathcal{C} \rightarrow M \rightarrow \mathbf{1}^r \rightarrow 0.$$

- We let

$$\Psi = \begin{bmatrix} \Omega & 0 & \cdots & 0 \\ \Omega L_{\alpha_1} & 1 & \cdots & 0 \\ \vdots & \vdots & \ddots & \vdots \\ \Omega L_{\alpha_r} & 0 & \cdots & 1 \end{bmatrix}.$$

Then

$$\Psi^{(-1)} = \Phi \Psi.$$

- Specialize Ψ at $t = \theta$ and find

$$\Psi(\theta) = \begin{bmatrix} -1/\pi_q & 0 & \cdots & 0 \\ -\log_C(\alpha_1)/\pi_q & 1 & \cdots & 0 \\ \vdots & \vdots & \ddots & \vdots \\ -\log_C(\alpha_r)/\pi_q & 0 & \cdots & 1 \end{bmatrix}.$$

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- Thus we can determine

$$\text{tr. deg}_{\bar{k}} \bar{k}(\pi_q, \log_C(\alpha_1), \dots, \log_C(\alpha_r))$$

by calculating

$$\dim \Gamma_\Psi.$$

Calculating Γ_Ψ

- Set $\Psi_1, \Psi_2 \in \mathrm{GL}_{r+1}(\mathbb{L} \otimes_{\bar{k}(t)} \mathbb{L})$ so that

$$(\Psi_1)_{ij} = \Psi_{ij} \otimes 1, \quad (\Psi_2)_{ij} = 1 \otimes \Psi_{ij},$$

and set $\tilde{\Psi} = \Psi_1^{-1} \Psi_2 \in \mathrm{GL}_{r+1}(\mathbb{L} \otimes_{\bar{k}(t)} \mathbb{L})$.

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- Define an $\mathbb{F}_q(t)$ -algebra map,

$$\mu = (X_{ij} \mapsto \tilde{\Psi}_{ij}) : \mathbb{F}_q(t)[X, 1/\det X] \rightarrow \mathbb{L} \otimes_{\bar{k}(t)} \mathbb{L},$$

which defines the $\mathbb{F}_q(t)$ -subgroup scheme $\Gamma_\Psi \subseteq \mathrm{GL}_{r+1}/\mathbb{F}_q(t)$.

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- In our case, this implies first that

$$\Gamma_\Psi \subseteq \left\{ \begin{bmatrix} * & 0 \\ * & \mathrm{id}_r \end{bmatrix} \right\} \subseteq \mathrm{GL}_{r+1}/\mathbb{F}_q(t).$$

Thus we can consider the coordinate ring of Γ_Ψ to be a quotient of $\mathbb{F}_q(t)[X_0, \dots, X_r, 1/X_0]$.

The vector group V

- The homomorphism of $\mathbb{F}_q(t)$ -group schemes

$$\begin{bmatrix} \alpha & 0 \\ \delta & \text{id}_r \end{bmatrix} \mapsto \alpha : \Gamma_\Psi \xrightarrow{\text{pr}} \mathbb{G}_m$$

coincides with the surjection,

$$\Gamma_\Psi \twoheadrightarrow \Gamma_C. \quad (\Gamma_C \cong \mathbb{G}_m).$$

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- Thus we have exact sequence of group schemes over $\mathbb{F}_q(t)$:

$$0 \rightarrow V \rightarrow \Gamma_\Psi \xrightarrow{\text{pr}} \mathbb{G}_m \rightarrow 0,$$

and we can consider $V \subseteq (\mathbb{G}_a)^r$ over $\mathbb{F}_q(t)$.

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- Now V is smooth over $\mathbb{F}_q(t)$ because $\text{pr} : \Gamma_\Psi \rightarrow \mathbb{G}_m$ is surjective on Lie algebras.
- It follows that defining equations for V are linear forms in X_1, \dots, X_r over $\mathbb{F}_q(t)$.

Defining equations for Γ_Ψ

- Pick $b_0 \in \mathbb{F}_q(t)^\times \setminus \mathbb{F}_q^\times$.
- Lift (use Hilbert Thm. 90) to

$$\gamma = \begin{bmatrix} b_0 & 0 & \cdots & 0 \\ b_1 & 1 & \cdots & 0 \\ \vdots & \vdots & \ddots & \vdots \\ b_r & 0 & \cdots & 1 \end{bmatrix} \in \Gamma_\Psi(\mathbb{F}_q(t)).$$

- We can use γ to create defining equations for Γ_Ψ using defining forms for V .

Theorem

- Suppose $F = c_1X_1 + \cdots + c_rX_r$, $c_1, \dots, c_r \in \mathbb{F}_q(t)$, is a defining linear form for V . Then

$$G = (b_0 - 1)F - F(b_1, \dots, b_r)(X_0 - 1)$$

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$$(b_0(\theta) - 1) \sum_{i=1}^r c_i(\theta) \log_C(\alpha_i) - \sum_{i=1}^r c_i(\theta) b_i(\theta) \pi_q = 0.$$

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- Let N be the k -linear span of $\pi_q, \log_C(\alpha_1), \dots, \log_C(\alpha_r)$. Then

$$\dim \Gamma_\Psi = \dim_k N.$$

Algebraic independence of Carlitz logarithms

- Starting with $\alpha_1, \dots, \alpha_r \in \bar{k}$ (suitably small), we found $\Phi \in \text{Mat}_r(\bar{k}[t])$ and $\Psi \in \text{Mat}_r(\mathbb{E})$ so that

$$\Psi(\theta) = \begin{bmatrix} -1/\pi_q & 0 & \cdots & 0 \\ -\log_C(\alpha_1)/\pi_q & 1 & \cdots & 0 \\ \vdots & \vdots & \ddots & \vdots \\ -\log_C(\alpha_r)/\pi_q & 0 & \cdots & 1 \end{bmatrix}.$$

- Since $\text{tr. deg}_{\mathbb{E}\bar{k}} \bar{k}(\pi_q, \log_C(\alpha_1), \dots, \log_C(\alpha_r)) = \dim \Gamma_\Psi = \dim_k N$, we can prove the following the following theorem.

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Theorem (P. 2008)

Suppose $\log_C(\alpha_1), \dots, \log_C(\alpha_r)$ are linearly independent over $k = \mathbb{F}_q(\theta)$. Then they are algebraically independent over \bar{k}

An Example

- Recall

$$\zeta_\theta = \sqrt[q]{-\theta}, \quad \exp_C(\pi_q/\theta) = \zeta_\theta, \quad \log_C(\zeta_\theta) = \frac{\pi q}{\theta}.$$

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$$\Phi = \begin{bmatrix} t - \theta & 0 \\ \zeta_\theta^{1/q}(t - \theta) & 1 \end{bmatrix}, \quad \Psi = \begin{bmatrix} \Omega & 0 \\ \Omega L_{\zeta_\theta} & 1 \end{bmatrix}.$$

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- We have a relation over \bar{k} on the entries of

$$\Psi(\theta) = \begin{bmatrix} -1/\pi_q & 0 \\ -1/\theta & 1 \end{bmatrix},$$

namely

$$\theta X_{21} + 1 = 0.$$

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- We begin with matrices in $\mathrm{GL}_r(\mathbb{L} \otimes_{\bar{k}(t)} \mathbb{L})$:

$$\psi_1 = \begin{bmatrix} \Omega \otimes 1 & 0 \\ \Omega L_{\zeta_\theta} \otimes 1 & 1 \end{bmatrix}, \quad \psi_2 = \begin{bmatrix} 1 \otimes \Omega & 0 \\ 1 \otimes \Omega L_{\zeta_\theta} & 1 \end{bmatrix}.$$

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Then the defining equations over $\mathbb{F}_q(t)$ for Γ_ψ will be precisely relations among the entries of

$$\psi_1^{-1} \psi_2 = \begin{bmatrix} \frac{1}{\Omega} \otimes \Omega & 0 \\ -L_{\zeta_\theta} \otimes \frac{1}{\Omega} + 1 \otimes \Omega L_{\zeta_\theta} & 1 \end{bmatrix}.$$

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- Consider the identity of functions (check!),

$$\zeta_\theta(t - \theta)\Omega(t) - t\Omega(t)L_{\zeta_\theta}(t) - 1 = 0,$$

and substitute into the lower left entry of $\psi_1^{-1} \psi_2$.

$$\psi_1^{-1}\psi_2 = \begin{bmatrix} \frac{1}{\Omega} \otimes \Omega & 0 \\ -L_{\zeta_\theta} \otimes \Omega + 1 \otimes \Omega L_{\zeta_\theta} & 1 \end{bmatrix}$$

$$\zeta_\theta(t - \theta)\Omega - t\Omega L_{\zeta_\theta} - 1 = 0$$

$$\Psi_1^{-1}\Psi_2 = \begin{bmatrix} \frac{1}{\Omega} \otimes \Omega & 0 \\ -L_{\zeta_\theta} \otimes \Omega + 1 \otimes \Omega L_{\zeta_\theta} & 1 \end{bmatrix}$$

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- Lower left entry of $\Psi_1^{-1}\Psi_2$ is

$$\begin{aligned} -L_{\zeta_\theta} \otimes \Omega + 1 \otimes \Omega L_{\zeta_\theta} &= -\left(\frac{1}{t}\zeta_\theta(t - \theta) - \frac{1}{t\Omega}\right) \otimes \Omega \\ &\quad + 1 \otimes \frac{1}{t}(\zeta_\theta(t - \theta)\Omega - 1) \\ &= -\frac{1}{t}\zeta_\theta(t - \theta) \otimes \Omega + \frac{1}{t\Omega} \otimes \Omega \\ &\quad + 1 \otimes \frac{1}{t}\zeta_\theta(t - \theta)\Omega - 1 \otimes \frac{1}{t} \\ &= \frac{1}{t}\left(\frac{1}{\Omega} \otimes \Omega\right) - \frac{1}{t}(1 \otimes 1). \end{aligned}$$

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- Therefore, Γ_Ψ is defined by

$$\Gamma_\Psi : tX_{12} - X_{11} + 1 = 0.$$

Carlitz zeta values

- Brief review of Carlitz zeta values
- Algebraic independence theorem of Chang-Yu
- Theorem of Chang-P.-Yu for varying q

Applications to Carlitz zeta values

$$\zeta_C(n) = \sum_{\substack{a \in \mathbb{F}_q[\theta] \\ a \text{ monic}}} \frac{1}{a^n} \in k_\infty, \quad n = 1, 2, \dots$$

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- As you may recall from the 2nd lecture, using the theory of Anderson and Thakur, one can construct a system of difference equations $\Psi^{(-1)} = \Phi\Psi$ so that $\zeta_C(n)$ appears in $\Psi(\theta)$.

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- As you may recall from the 2nd lecture, using the theory of Anderson and Thakur, one can construct a system of difference equations $\Psi^{(-1)} = \Phi\Psi$ so that $\zeta_C(n)$ appears in $\Psi(\theta)$.
- Known algebraic relations over \bar{k} among $\zeta_C(n)$:

$$(q-1) \mid n \Rightarrow \zeta_C(n) = r_n \pi_q^n, \quad r_n \in \mathbb{F}_q(\theta), \quad (\text{Euler-Carlitz})$$

$$\zeta_C(np) = \zeta_C(n)^p, \quad (\text{Frobenius}).$$

The Chang-Yu Theorem

Algebraic independence of $\zeta_C(n)$

Theorem (Chang-Yu 2007)

For any positive integer n , the transcendence degree of the field

$$\bar{k}(\pi_q, \zeta_C(1), \dots, \zeta_C(n))$$

over \bar{k} is

$$n - \left\lfloor \frac{n}{p} \right\rfloor - \left\lfloor \frac{n}{q-1} \right\rfloor + \left\lfloor \frac{n}{p(q-1)} \right\rfloor + 1.$$

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Question: What can we say about Carlitz zeta values if we allow q to vary?

The Chang-Yu Theorem

Algebraic independence of $\zeta_C(n)$

Theorem (Chang-Yu 2007)

For any positive integer n , the transcendence degree of the field

$$\bar{k}(\pi_q, \zeta_C(1), \dots, \zeta_C(n))$$

over \bar{k} is

$$n - \left\lfloor \frac{n}{p} \right\rfloor - \left\lfloor \frac{n}{q-1} \right\rfloor + \left\lfloor \frac{n}{p(q-1)} \right\rfloor + 1.$$

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Answer: Even then, the Euler-Carlitz relations and the Frobenius p -th power relations tell the whole story....

Zeta values with varying constant fields

For $m \geq 1$, we set

$$\zeta_m(n) = \sum_{\substack{a \in \mathbb{F}_{p^m}[\theta] \\ a \text{ monic}}} \frac{1}{a^n}, \quad n = 1, 2, \dots$$

Theorem (Chang-P.-Yu)

For any positive integers s and d , the transcendence degree of the field

$$\bar{k}(\cup_{m=1}^d \{\pi p^m, \zeta_m(1), \dots, \zeta_m(s)\})$$

over \bar{k} is

$$\sum_{m=1}^d \left(s - \left\lfloor \frac{s}{p} \right\rfloor - \left\lfloor \frac{s}{p^m - 1} \right\rfloor + \left\lfloor \frac{s}{p(p^m - 1)} \right\rfloor + 1 \right).$$

Rank 2 Drinfeld modules

- Periods and quasi-periods
- A Galois group example
- Algebraic independence in the non-CM case

Periods and quasi-periods of rank 2 Drinfeld modules

- Recall that for a rank 2 Drinfeld module $\rho : \mathbb{F}_q[t] \rightarrow \bar{k}[F]$ with

$$\rho(t) = \theta + \kappa F + F^2,$$

we can take

$$\Phi = \begin{bmatrix} 0 & 1 \\ t - \theta & -\kappa^{1/q} \end{bmatrix}, \quad \Psi = \begin{bmatrix} 0 & 1 \\ 1 & -\kappa \end{bmatrix} \begin{bmatrix} s_1^{(1)} & s_1^{(2)} \\ s_2^{(1)} & s_2^{(2)} \end{bmatrix}^{-1}.$$

- Furthermore,

$$\Psi(\theta)^{-1} = \begin{bmatrix} \omega_1 & \eta_1 \\ \omega_2 & \eta_2 \end{bmatrix},$$

where $\omega_1, \omega_2, \eta_1, \eta_2$ are the periods and quasi-periods for ρ .

An example ($\kappa = \sqrt{\theta} + \sqrt{\theta^q}$)

- Assume $p \neq 2$. Consider the Drinfeld module ρ with

$$\rho(t) = \theta + (\sqrt{\theta} + \sqrt{\theta^q})F + F^2.$$

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$$\Gamma_\Psi = \left\{ \begin{bmatrix} \alpha & \beta t \\ \beta & \alpha \end{bmatrix} \right\}.$$

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Theorem (Thiery 1992)

The period matrix of a Drinfeld module of rank 2 over \bar{k} with CM has transcendence degree 2 over \bar{k} .

Rank 2 Drinfeld modules without CM

In general, we say that a Drinfeld module ρ *does not have complex multiplication* if

$$\text{End}(\rho) = \mathbb{F}_q[t].$$

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Theorem (Chang-P.)

Suppose that $p \neq 2$. Let ρ be a Drinfeld module of rank 2 over \bar{k} without CM. Then

$$\Gamma_\rho \cong \text{GL}_2.$$

In particular, the periods and quasi-periods of ρ ,

$$\omega_1, \omega_2, \eta_1, \eta_2,$$

are algebraically independent over \bar{k} .